

# Phase separation and metal-insulator transitions in the spin-one-half Falicov-Kimball model

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## Abstract

The ground-state phase diagram of the spin-one-half Falicov-Kimball model (FKM) is studied in the one dimension using small-cluster exact-diagonalization calculations. The resultant exact solutions are used to examine possibilities for valence and metal-insulator transitions in this model. A number of remarkable results are found. (i) The phase separation in the spin-one-half FKM takes place for a wide range of  $f$ -electron concentrations  $n_f$  and  $d-f$  interactions  $G$ , including  $G$  large. (ii) In the strong coupling limit ( $G > 4$ ) the model exhibits a pressure induced discontinuous insulator-metal transition from an integer-valence state ( $n_f = 1$ ) into another integer-valence state ( $n_f = 0$ ). (iii) For intermediate values of  $G$  ( $G \sim 2.5$ ) the FKM undergoes a few discontinuous intermediate-valence transitions. There are several discontinuous insulator-insulator transitions from  $n_f = 1$  to  $n_f = 1/2$  and a discontinuous insulator-metal transition from  $n_f = 1/2$  to  $n_f = 0$ . (iv) In the weak coupling limit ( $G < 2$ ) the model undergoes a few consecutive discontinuous and continuous intermediate-valence transitions as well as a discontinuous metal-insulator transition.

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## 1 Introduction

Recent new measurements of the insulator-metal transitions in SmB<sub>6</sub> [1] and in the transition-metal halides [2], have once again sparked an interest in valence and insulator-metal transitions [3]. These transitions are observed in a wide group of substances formed by transition-metal oxides as well as rare-earth sulfides and borides, when some external parameters (like pressure or temperature) are varied. They are in many cases first-order phase transitions, however, second-order transitions ranging from very gradual to rather steep are also observed [4].

To describe all such transitions in a unified picture Falicov and Kimball [5] introduced a simple model in which only two relevant single-electron states are taken into account: extended Bloch waves and a set of localized states centered at the sites of the metallic ions in the crystal. It is assumed that insulator-metal transitions result from a change in the occupation numbers of these electronic states, which remain themselves basically unchanged in their character. The Hamiltonian of the model can be written as the sum of four terms:

$$H_0 = \sum_{ij\sigma} t_{ij} d_{i\sigma}^+ d_{j\sigma} + G \sum_{i\sigma\sigma'} f_{i\sigma}^+ f_{i\sigma} d_{i\sigma'}^+ d_{i\sigma'} + E_f \sum_{i\sigma} f_{i\sigma}^+ f_{i\sigma} + \frac{U}{2} \sum_{i\sigma} f_{i\sigma}^+ f_{i\sigma} f_{i-\sigma}^+ f_{i-\sigma}, \quad (1)$$

where  $f_{i\sigma}^+$ ,  $f_{i\sigma}$  are the creation and annihilation operators for an electron of spin  $\sigma$  in the localized state at lattice site  $i$  with binding energy  $E_f$  and  $d_{i\sigma}^+$ ,  $d_{i\sigma}$  are the creation and annihilation operators of the itinerant electrons in the  $d$ -band Wannier state at site  $i$ .

The first term of (1) is the kinetic energy corresponding to quantum-mechanical hopping of the itinerant  $d$  electrons between sites  $i$  and  $j$ . These intersite hopping transitions are described by the matrix elements  $t_{ij}$ , which are  $-t$  if  $i$  and  $j$  are the nearest neighbors and zero otherwise (in the following all parameters are measured in units of  $t$ ). The second term represents the on-site Coulomb interaction between the  $d$ -band electrons with density  $n_d = N_d/L = \frac{1}{L} \sum_{i\sigma} d_{i\sigma}^+ d_{i\sigma}$  and the localized  $f$  electrons with density  $n_f = N_f/L = \frac{1}{L} \sum_{i\sigma} f_{i\sigma}^+ f_{i\sigma}$ , where  $L$  is the number of lattice sites. The third term stands for the localized  $f$  electrons whose sharp energy level is  $E_f$ . The last term represents the intra-atomic Coulomb interaction between the localized  $f$  electrons.

In spite of the fact that the FKM is one of the simplest examples of interacting fermionic system, the theoretical picture of valence and insulator-metal transitions remains still uncertain in the framework of this model. Even, in the existing literature on this model, different answers can be found on the fundamental question whether the FKM can describe both the discontinuous and continuous changes of the  $f$  (d)-electron occupation number  $n_f$  ( $n_d$ ) as a function of the  $f$ -level energy  $E_f$  [6]. It should be noted that this question is indeed crucial for the systems mentioned above, since, supposing [4] that the external pressure shifts the energy level  $E_f$ , the valence changes observed in some rare-earth and transition-metal compounds (SmS, SmB<sub>6</sub>, Ti<sub>2</sub>O<sub>3</sub>, and so on) could be understandable purely electronically. Unfortunately, it was found that valence and insulator-metal transitions are very sensitive to the approximation used. Various approximations [6] (mean-field, virtual crystal, CPA, etc.) yield very different and often fully controversial results. This indicates that the study of valence and insulator-metal transitions may be successful only with methods which are relatively insensitive to the type of approximation used and, of course, with exact methods.

In our previous papers [7, 8] we have showed that the method of small-cluster exact-diagonalization calculations is very effective in describing ground-state properties of both the spinless and spin-one-half FKM. For the spinless FKM [7] the exact numerical calculations (over the full set of  $f$ -electron configurations) can be performed on relatively large clusters ( $L \sim 36$ ), even without some special computational effort. For such clusters the finite-size effects are practically negligible and the results can be satisfactorily extrapolated to the thermodynamics limit ( $L \rightarrow \infty$ ). Using this method we have described successfully the ground-state phase diagram as well as the picture of valence and metal-insulator transitions in the spinless FKM [7]. The situation for the spin-one-half FKM is more complicated. The full Hilbert space of the spin-one-half FKM is much larger than one of the spinless FKM and this fact impose severe restrictions on the size of clusters that can be studied by the exact-diagonalization method. Our recent [8] numerical computations performed for the spin-one-half FKM with  $U$  and  $G$  finite showed that clusters with  $L > 24$  are beyond the reach of present day computers. Fortunately, the size of the Hilbert space can be reduced considerably in some special, but physically still interesting limits, e.g.,  $U \rightarrow \infty$ . In the limit  $U \rightarrow \infty$  states with two  $f$  electrons at the same site are projected out thereby much larger clusters ( $L \sim 36$ ) become accessible for the numerical investigation in

this reduced subspace. For this reason all calculations presented in this paper have been done at  $U = \infty$ . The main goal for performing these calculations was to construct the comprehensive phase diagram of the spin-one-half FKM. The second goal of our numerical study was to find and describe all possible types of valence and metal-insulator transitions in this model. We show that both the ground-state phase diagram as well as the picture of valence and metal-insulator transitions obtained for  $U = \infty$  strongly differ from ones obtained in our previous paper [8] for a restricted set of  $f$ -electron configurations and  $U$  finite. In particular, the exhaustive numerical studies of the model (at  $U = \infty$ ) performed on finite clusters up to 36 sites revealed some new unexpected features like the phase separation for all nonzero  $G$  and discontinuous metal-insulator transitions for  $G$  small, to mention only a few. In the present paper we discuss these features in detail.

Since the  $f$ -electron density operators  $f_{i\sigma}^+ f_{i\sigma}$  of each site  $i$  commute with the Hamiltonian (1), the  $f$ -electron occupation number is a good quantum number, taking only two values,  $w_i = 0, 1$  according to whether the site  $i$  is unoccupied or occupied by the localized  $f$  electron (configurations with  $w_i = 2$  are projected out due to  $U = \infty$ ). Therefore the Hamiltonian (1) can be replaced by

$$H = \sum_{ij\sigma} h_{ij} d_{i\sigma}^+ d_{j\sigma} + E_f \sum_i w_i, \quad (2)$$

where  $h_{i,j} = t_{ij} + Gw_i\delta_{ij}$ .

Thus for a given  $f$ -electron configuration  $w = \{w_1, w_2 \dots w_L\}$  defined on a one-dimensional lattice with periodic boundary conditions, the Hamiltonian (2) is the second-quantized version of the single-particle Hamiltonian  $h(w) = T + GW$ , so the investigation of the model (2) is reduced to the investigation of the spectrum of  $h$  for different configurations of  $f$  electrons. Since the  $d$  electrons do not interact among themselves, the numerical calculations precede directly in the following steps (we consider only the case  $N_f + N_d = L$ , which is the point of the special interest for valence and metal-insulator transitions caused by promotion of electrons from localized  $f$  orbitals ( $f^n \rightarrow f^{n-1}$ ) to the conduction band states). (i) Having  $G$ ,  $E_f$  and  $w = \{w_1, w_2 \dots w_L\}$  fixed, find all eigenvalues  $\lambda_k$  of  $h(w) = T + GW$ . (ii) For a given  $N_f = \sum_i w_i$  determine the ground-state energy  $E(w) = \sum_{k=1}^{L-N_f} \lambda_k + E_f N_f$  of a particular  $f$ -electron configuration  $w$  by filling in the lowest  $N_d = L - N_f$  one-electron levels (the spin degeneracy must be taken into account). (iii) Find the  $w^0$  for which  $E(w)$  has a minimum. Repeating this procedure for different

values of  $N_f$ ,  $G$  and  $E_f$ , one can immediately study the ground-state phase-diagram of the model in the  $N_f$ - $G$  plane as well as the dependence of the  $f$ -electron occupation number  $N_f = \sum_i w_i^0$  on the  $f$ -level position  $E_f$  (valence transitions).

## 2 Phase separation

Let us first discuss the ground-state phase-diagram of the spin-one-half FKM. To reveal the basic structure of the phase diagram in the  $N_f$ - $G$  plane ( $E_f = 0$ ) we have performed an exhaustive study of the model on finite (even) clusters up to 36 sites. For fixed  $L$  the numerical calculations have been done along the lines discussed above with a step  $\Delta N_f = 2$  and  $\Delta G = 0.05$ . The results of numerical computations are summarized in Fig. 1. These results show that the phase diagram of the spin-one-half FKM consists of three main domains: the most homogeneous domain MHD (in Fig. 1 denoted as  $\cdot$ ) and two phase separation domains  $PSD_1$  (denoted as  $\circ$ ,  $\times$  and  $+$ ) and  $PSD_2$  (denoted as  $\triangle$ ). In the MHD the ground states are configurations in which the atomic or  $n$ -molecule clusters of  $f$ -electrons are distributed in such a manner that the distances between two consecutive clusters are either  $d$  or  $d + 2$ . Furthermore the distribution of the distances of  $d$  and  $d + 2$  has to be most homogeneous. Two basic types of the ground state configurations that fill up practically the whole MHD are displayed in Table. 1. In the  $PSD_1$  the ground state configurations are mixtures  $w \& w_e$  of aperiodic (periodic) configurations  $w$  and the empty configuration  $w_e = \{00\dots 0\}$  (see Ref. [9] for a definition of mixtures), i.e., all  $f$ -electrons are distributed only in one part of the lattice (configuration  $w$ ) while another (connecting) part is free of  $f$ -electrons (phase separation). We have found three basic types of configurations  $w$  which form these mixtures: (i) the aperiodic atomic configurations  $w_a = \{10_{k_1}10_{k_2}\dots 10_{k_i}1\dots 0_{k_2}10_{k_1}1\}$  (with  $k_i > 0$ ), (ii) the aperiodic  $n_i$ -molecule configurations  $w_b = \{1_{n_1}0_{k_1}1_{n_2}0_{k_2}\dots 1_{n_i}0_{k_i}1_{n_1}\dots 0_{k_2}1_{n_2}0_{k_1}1_{n_1}\}$  (with  $1 \leq n_i \leq N_f/2$  and  $k_i > 0$ ), and (iii) the  $N_f$ -molecule (segregated) configurations  $w_c = \{11\dots 1\}$ . Three different regions of stability corresponding to mixtures  $w_a \& w_e$ ,  $w_b \& w_e$  and  $w_c \& w_e$  are denoted in Fig. 1 as  $\circ$ ,  $\times$  and  $+$ . It is seen that the mixtures of the atomic configuration  $w_a$  and the empty configuration are stable only at low  $f$ -electron concentrations and the Coulomb interactions  $G < 2.2$ . A direct comparison of results obtained for the spin-one-half and spinless FKM [9] shows that this region corresponds roughly to a region of phase separation in the spinless FKM. Outside these regions the phase diagrams of the spin-one-half and spinless FKM are,

however, strongly different. While the phase separation in the spinless FKM takes place only for weak interactions ( $G < 1.2$ ) the spin-one-half FKM exhibits the phase separation for all Coulomb interactions. Even with increasing  $G$  the phase separation shifts to higher  $f$ -electron concentrations. Particularly, in the region  $2.5 < G < 2.7$  where the ground states are the mixtures of the  $n$ -molecule configurations  $w_b$  with the empty configuration the phase separation takes place for all  $N_f < L/2$  and in the region  $G > 2.7$  where the ground states are the segregated configurations  $w_S = w_c \& w_e$  even for all  $N_f < L$ . At large  $f$ -electron concentrations  $n_f$  but in the opposite limit ( $G < 0.4$ ) there exists another small domain of the phase separation  $PSD_2$  (denoted as  $\Delta$ ). The numerical results on finite lattices up to 36 sites revealed only one type of configurations that can be the ground state configurations in this domain, and namely, the mixtures of the periodic  $n$ -molecule configurations  $w_d$  with the fully occupied lattice  $w_f = \{11\dots1\}$  (the length of a connected cluster of occupied sites in these mixtures is at least  $L/2$ ).

The second step in our numerical studies has been the extrapolation of small-cluster exact diagonalization results on large lattices. In Fig. 2 we present the ground state phase diagram of the spin-one-half FKM obtained for  $L = 100$  on the extrapolated set of configurations that includes practically all possible types of the ground-state configurations found on finite lattices up to 36 sites. In particular, we have considered (i) the most homogeneous configurations  $w_h$  of the type  $a$  and  $b$  (see Table 1), (ii) all mixtures  $w_a \& w_e$  with  $k_i$  smaller than 6, (iii) all mixtures  $w_b \& w_e$  with  $n_i$  and  $k_i$  smaller than 6, (iv) all mixtures  $w_d \& w_f$  with periods smaller than 12, and (v) all segregated configurations. One can see that all fundamental features of the phase diagram found on small lattices hold on much larger lattices too. Of course, the phase boundaries of different regions corresponding to mixtures  $w_a \& w_e$  ( $a$ ),  $w_b \& w_e$  ( $b$ ),  $w_c \& w_e$  ( $c$ ) and  $w_d \& w_f$  ( $d$ ) are now more obvious. Since the mixtures  $w_a \& w_e$ ,  $w_b \& w_e$ ,  $w_c \& w_e$  and  $w_d \& w_f$  are metallic [9] one can expect that the phase boundary between the MHD and PSD is also the boundary of the correlation induced metal-insulator transition. To confirm this conjecture it is necessary to show that the most homogeneous configurations from the MHD are insulating, or by other words that there is a finite energy gap  $\Delta$  at the Fermi energy in the spectra of these configurations [10]. The numerical results for the most-homogeneous configurations of the type  $a$  and  $b$  that fill up practically the whole MHD are displayed in Fig. 3. The insulating character of these configurations is clearly demonstrated by the finite  $\Delta$  that exists for all nonzero  $f$ -electron

concentrations  $n_f$  and Coulomb interactions  $G$ . Thus we can conclude that the spin-one-half FKM undergoes (on the phase boundary between the MHD and PSD) the correlation induced metal-insulator transition that is accompanied by a discontinuous change of the energy gap  $\Delta$ .

### 3 Valence and metal-insulator transitions

The existence of a large metallic domain in the strong coupling region is the main difference between the phase diagram of the spin-one-half and spinless FKM. In the spinless FKM [9] the existence of the metallic phase was restricted only on a small region  $G < 1.2$  and  $n_f < 1/4$  ( $n_f > 3/4$ ), while the remaining part of the phase diagram was insulating. The phase diagram of the spin-one-half FKM has a more complicated structure. In addition to the insulating phase corresponding to the homogeneous configurations  $w_h$  there are four metallic phases corresponding to four different classes of the ground state configurations:  $w_a \& w_e$ ,  $w_b \& w_e$ ,  $w_c \& w_e$  and  $w_d \& w_f$ . Of course, this fact has to lead also to different picture of valence and metal-insulator transitions induced by pressure (increasing  $E_f$ ). Let us now discuss, in more detail, possible types of these transitions. In order to examine possible types of valence and insulator-metal transitions induced by pressure we have chose three different values of  $G$  ( $G = 1, 2.5, 5$ ) that represent three typical regimes of the spin-one-half phase diagram. For  $G = 5$  the valence transition can be constructed immediately. Indeed the ground states in this region are only the segregated configurations and thus they can be used directly in numerical calculations instead of the full set of  $f$ -electron configurations. This allows to perform numerical calculations on large systems and practically fully exclude finite size effects. Results of numerical calculations are shown in Fig. 4. It is seen that in the strong coupling limit the spin-one-half FKM exhibits a discontinuous valence transition from an integer valence ground state  $n_f = 1$  into another integer valence ground state  $n_f = 0$  at  $E_f = E_c \sim 1.273$ . The inset in Fig. 4 shows that the discontinuous valence transition is accompanied by a discontinuous insulator-metal transition since the energy gap  $\Delta$  vanishes discontinuously at  $E_f = E_c$ . Thus we can conclude that the spin-one-half FKM in the pressure induced case can describe the discontinuous insulator-metal transitions from an integer-valence state ( $n_f = 1$ ) into another integer-valence state ( $n_f = 0$ ).

The situation for  $G = 2.5$  is slightly complicated. Although we know that the ground states for  $2.45 < G < 2.75$  (and  $n_f < 1/2$ ) are configurations of a type  $w_b \& w_e$  the number

of configurations belonging to this class is still too large for numerical calculations on large lattices and thus it should be further reduced. For this reason we have performed numerical calculations for all finite (even) clusters up to 48 sites at selected value  $G = 2.5$ . We have found that for any  $N_f > 0$  only one type of configurations and namely mixtures  $\{11001100 \dots 1100\} \& w_e$  are the ground states at  $G = 2.5$ . This allows us to avoid technical difficulties associated with a large number of configurations and consequently to study very large systems ( $L \sim 1200$ ). The valence transition obtained for this set of configurations (and  $w_h$  configurations for  $n_f > 1/2$ ) is shown in Fig. 5. It is seen that the valence transition for  $G = 2.5$  is very steep. A more detail analysis showed, however, that the transition from  $n_f = 1$  to  $n_f = 0$  is not a discontinuous integer valence transition but it consists of several discontinuous intermediate valence transitions. Particularly, there are several discontinuous insulator-insulator transitions from  $n_f = 1$  to  $n_f = 1/2$  and a discontinuous insulator-metal transition from  $n_f = 1/2$  to  $n_f = 0$ . To exclude the possibility that this structure is a consequence of the finite size of clusters used in numerical calculations we have performed calculations on several large clusters ( $L = 400, 800, 1200$ ) but no finite size effects on the valence transition were observed.

The most complicated situation is for small and intermediate values of  $G$ . For example, at selected value of  $G = 1$  the ground states are the insulating configurations  $w_h$  (for  $n_f > 0.16$ ) and the metallic configurations  $w_a \& w_e$  (for  $n_f < 0.16$ ). However, for a construction of the valence transition on large lattices this knowledge is too general and thus for  $n_f < n_c$  it is again necessary to perform an additional study of the model in order to determine the explicit type of ground-state configurations at selected value of  $G$ . The exhaustive numerical studies that we have performed for  $n_f < n_c$  on finite lattices up to 60 sites showed that only configurations of the type  $w_1 = \{1010_{n_1} 1010_{n_2} \dots 0_{n_2} 1010_{n_1} 101\} \& w_e$  or configurations of the type  $w_2 = \{10_{n_1} 1010_{n_2} 1010_{n_3} \dots 0_{n_3} 1010_{n_2} 1010_{n_1} 1\} \& w_e$  can be the ground states at  $G = 1$ . This fact allows us to use the set of configurations  $w_1$  and  $w_2$  instead of the much larger set  $w_a \& w_e$  and consequently to construct the valence transition on large lattices ( $L \sim 1200$ ). The results of numerical calculations obtained for the  $E_f$ -dependence of the  $f$ -electron occupation number  $n_f$  and the energy gap  $\Delta$  are displayed in Fig. 6. Unlike the valence transitions obtained for  $G = 5$  and  $G = 2.5$  the valence transitions in the weak coupling limit are much wider and consist of several consecutive discontinuous and continuous valence transitions. The valence transitions for

$E_f < E_c \sim -1.05$  are insulator-insulator transitions since they realize between insulating ground states corresponding to the most homogeneous configurations  $w_h$ . At  $E_f = E_c$  the spin-one-half FKM undergoes a pressure induced discontinuous valence transition from an intermediate valence state with  $n_f \sim 0.19$  into another intermediate-valence state with  $n_f \sim 0.14$  that is accompanied by a discontinuous insulator-metal transition. Above  $E_c$  the  $f$ -electron occupation number  $n_f$  changes continuously and vanishes at  $E_f = E_0 \sim -0.8$ .

It should be noted that the weak-coupling picture of valence and metal-insulator transitions presented in this paper strongly differs from one found in our previous work [8] for the restricted set of configurations consisting of only the most homogeneous configurations  $w_h$ . For such a set of configurations the system exhibits the metal-insulator transition only at  $E_f = E_0$  and this transition is continuous. This fact indicated the serious deficiency of the spin-one-half FKM since in some rare-earth compounds (e.g. in SmB<sub>6</sub> [1]) just a discontinuous insulator-metal transition is observed. Fortunately, a more accurate analysis performed in the present paper has not confirmed this deficiency of the model. Contrary, our numerical results showed that the spin-one-half FKM is capable of describing not only a discontinuous metal-insulator transition but all basic types of valence transitions observed experimentally in rare-earth and transitions-metal compounds. This indicates that the spin-one-half FKM could yield the correct physics for describing these materials.

In summary, the extrapolation of small-cluster exact-diagonalization calculations has been used to examine the ground-state phase diagram of the spin-one-half FKM in the one dimension. A number of remarkable results have been found. (i) The phase separation in the spin-one-half FKM takes place for a wide range of  $f$ -electron concentrations  $n_f$  and  $d-f$  interactions  $G$ , including  $G$  large. (ii) In the strong coupling limit ( $G > 4$ ) the model exhibits a pressure induced discontinuous insulator-metal transition from an integer-valence state ( $n_f = 1$ ) into another integer-valence state ( $n_f = 0$ ). (iii) For intermediate values of  $G$  ( $G \sim 2.5$ ) the FKM undergoes a few discontinuous intermediate-valence transitions. There are several discontinuous insulator-insulator transitions from  $n_f = 1$  to  $n_f = 1/2$  and a discontinuous insulator-metal transition from  $n_f = 1/2$  to  $n_f = 0$ . (iv) In the weak coupling limit ( $G < 2$ ) the model undergoes a few consecutive discontinuous and continuous intermediate-valence transitions as well as a discontinuous metal-insulator transition.

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- [10] Since we consider the case  $N_f + N_d = L$ , the Fermi level  $E_F$  and the energy gap  $\Delta$  of some configuration  $w$  are given by  $E_F = \lambda_{L-N_f}$  and  $\Delta = \lambda_{L-N_f+1} - \lambda_{L-N_f}$ , respectively.

## Figure Captions

Fig. 1. The ground-state phase diagram of the spin-one-half FKM obtained over the full set of  $f$ -electron configurations. For  $1/4 < n_f < 3/4$  numerical calculations have been done on the lattice with  $L = 28$  while for  $n_f \leq 1/4$  and  $n_f \geq 3/4$  on the lattice with  $L = 36$ . Four different regions of stability corresponding to mixtures  $w_a \& w_e$ ,  $w_b \& w_e$ ,  $w_c \& w_e$  and  $w_d \& w_f$  are denoted as  $\circ$ ,  $\times$ ,  $+$  and  $\Delta$ .

Fig. 2. The ground-state phase diagram of the spin-one-half FKM obtained on the extrapolated set of  $f$ -electron configurations for  $L = 100$ . Four different regions of stability corresponding to mixtures  $w_a \& w_e$ ,  $w_b \& w_e$ ,  $w_c \& w_e$  and  $w_d \& w_f$  are denoted as  $a, b, c$  and  $d$ . The dashed line through the exact numerical points is a guide to the eye.

Fig. 3.  $n_f$ -dependence of the energy gap  $\Delta$  calculated for the most homogeneous configurations of the type  $a$  and  $b$  (see Table 1).

Fig. 4. Dependence of the  $f$ -electron occupation number  $n_f$  on the  $f$ -level position  $E_f$  for  $G = 5$  and three different values of  $L$ . Inset: The behavior of  $n_f$  and  $\Delta$  close to the insulator-metal transition point.

Fig. 5. Dependence of the  $f$ -electron occupation number  $n_f$  on the  $f$ -level position  $E_f$  for  $G = 2.5$  and three different values of  $L$ . Inset: The behavior of  $n_f$  and  $\Delta$  close to the insulator-metal transition point.

Fig. 6. Dependence of the  $f$ -electron occupation number  $n_f$  on the  $f$ -level position  $E_f$  for  $G = 1$  and three different values of  $L$ . Inset: The behavior of  $n_f$  and  $\Delta$  close to the insulator-metal transition point.

### Table Captions

Table 1: Two basic types of the most homogeneous configurations that fill up practically the whole MHD for  $L = 24$ .

$N_f$	$w_h^a$	$w_h^b$
4	100001000000100001000000	110000000000001100000000000
6	100100001001000010010000	110000001100000011000000
8	100100100100100100100100	110000110000110000110000
10	110010011001001001100100	110011000011001100110000
12	110011001100110011001100	110011001100110011001100
14	111001110011100111001100	
16	111100111100111100111100	
18	111111001111110011111100	
20	111111111100111111111100	











